Exploring the astrophysical conditions for the creation of the first r–process peak, and the impact of nuclear physics uncertainties

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Abstract
We present a study of nucleosynthesis for conditions of high Ye outflows from Neutron Star Mergers (BNSs). We investigate the effect of new beta–decay rates measurements and uncertainties in nuclear masses of the newly measured \textsuperscript{84,85}Ga to the r–process nucleosynthesis calculations. The impact of these quantities to the production of the elements of the r–process abundance pattern for \textit{A<100} is quantified and presented. This proceedings paper is based on [1].

Keywords
r–process, nucleosynthesis

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INTRODUCTION

The rapid neutron capture process (r–process) is responsible for the production of about half of the heavy elements observed in the solar abundances. It is characterized by three peaks, each corresponding to the closed neutron shells at \textit{N} = 50, 82, 126. The peaks are the result of matter accumulating at the neutron closed shells. The drop in binding energy that characterizes the region after the shell closure leads to a bending of the path towards stability, reaching nuclei with longer β–decay half-lives and producing an accumulation of material, compared to nuclei before or beyond the shell closures.

A site of the r–process was unknown until recent observations. The gravitational wave event GW170817 \textsuperscript{[2,3]} which was identified as a Binary Neutron Star merger (BNSs) is considered as a site for the r–process nucleosynthesis \textsuperscript{[4–9]}. The detection of gravitational waves from the binary neutron star was followed by the detection of fast fading optical/infrared counterpart (AT2017gfo) \textsuperscript{[10]}, consistent with the predictions for a kilonova, associated with r–process nucleosynthesis \textsuperscript{[11–15]}, establishing the production of heavy elements in the aftermath of BNS \textsuperscript{[15–19]}. The complicated atomic structure of lanthanides implies high opacity ejecta; therefore, lanthanides emit light in the red wavelengths, thus, the blue color of the emission spectra at early times indicates that part of the ejecta is characterized by relatively high electron fraction (\textit{Y_e}) (0.25–0.4) and consequently low lanthanide production \textsuperscript{[20–26]}.

Assuming that BNSs are partially responsible for the production of first r–process peak elements, we present a study of nucleosynthesis under moderate \textit{Y_e} (0.35–0.38) and low entropy (S=10
kB/baryon) conditions. The conditions are consistent with the blue part of the kilonova observation (no lanthanides are created) and provide the strongest contribution to the mass region around the A=80 and A=84. Under these conditions, we investigate the effects of Ga isotopes mass uncertainties and the influence of recently published β–delayed neutron probabilities on the final r–process abundance pattern. We note that the range of $Y_e$ is justified because as shown in [1] lower $Y_e$ (<0.35) leads to overproduction of A=90–120 region whereas $Y_e$ above 0.39 leads to underproduction of the first r–process peak elements.

**METHODS**

We use the Hauser Feshbach statistical code TALYS [27] to calculate neutron capture reaction rates based on the mass values of the recently measured $^{84,85}$Ga. To systematically study the impact of $^{84,85}$Ga masses, on the formation of $A \approx 84$ nuclei we use a Monte–Carlo approach. We assume the "true" mass value is distributed following a normal distribution with 3σ according to the uncertainty, of their extrapolated mass values given in the atomic mass evaluation data (AME_16) [35] (200, 300 KeV respectively). Table 1 summarizes the new mass measurements and AME_16 corresponding values.

**Table 1.** Mass measurements of $^{84,85}$Ga isotopes reported in [1] in comparison to AME_16 [35]. The symbol # indicates systematics.

<table>
<thead>
<tr>
<th></th>
<th>Mass excess TITAN (keV/c²)</th>
<th>Mass excess AME_16 (keV/c²)</th>
<th>Difference (keV/c²)</th>
</tr>
</thead>
<tbody>
<tr>
<td>$^{84}$Ga</td>
<td>−44 094 (30)</td>
<td>−44 090 (200)#</td>
<td>4</td>
</tr>
<tr>
<td>$^{85}$Ga</td>
<td>−39 744 (37)</td>
<td>−39 850 (300)#</td>
<td>−106</td>
</tr>
</tbody>
</table>

Nuclear reaction rates not affected by the new Ga masses measurements were taken from JINA REACLIB [29]. Experimental masses from AME_16 were used when available; otherwise, we use the FRDM mass model [30]. We then use each set of the resulting neutron capture rates in GSINet [28] to calculate the r–process abundances.

β–decay rates and β–delayed neutron emission branches were taken from experimentally known nuclear properties database (NUBASE16) [31] when available. When experimental values were not available, values from theoretical predictions [32] were used. Here we included recent β–delayed neutron emission ($P_{1n}$) of $^{82,83,84,85}$Ga [33,34] to study the effect β–delayed–neutron emission branches to the final r–process abundance pattern. The new $P_{1n}$ measurements compared to NUBASE16 are summarized in Table 2.

**Table 2.** $P_{1n}$ of $^{82,83,84,85}$Ga isotopes reported in [33,34], in comparison to NUBASE16 [35]. The symbol # indicates systematics.

<table>
<thead>
<tr>
<th></th>
<th>Reported at [33,34] $P_{1n}$ (%)</th>
<th>NUBASE16 [31] $P_{1n}$ (%)</th>
<th>Difference (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>$^{82}$Ga</td>
<td>22 (2)</td>
<td>21.3 (13)</td>
<td>0.7</td>
</tr>
<tr>
<td>$^{83}$Ga</td>
<td>85 (4)</td>
<td>62.8 (25)</td>
<td>22.2</td>
</tr>
<tr>
<td>$^{84}$Ga</td>
<td>53 (20)</td>
<td>40 (7)</td>
<td>13</td>
</tr>
<tr>
<td>$^{85}$Ga</td>
<td>70 (5)</td>
<td>&gt;35 #</td>
<td>35</td>
</tr>
</tbody>
</table>

The thermodynamic evolution of the systems was parametrized assuming homologous expansion following [36]. The initial temperature was set at $T = 6$ GK, expansion timescale at 7 ms and entropy
at 10 kB/baryon. We perform calculations for $Y_e = 0.35–0.38$. Different electron fraction results were considered with equal weight.

**RESULTS**

We compare our results with the solar r–process abundances (Fig. 1) in the region $A=80–90$. The pink uncertainty bands show the variation of the abundances that arise from the uncertainties of the masses of $^{84,85}$Ga from AME$_{16}$. The new Ga mass values affect the abundances of elements with mass number $A=82$–$87$, with the biggest impact on $A=83$ accounting to $\approx 10\%$ despite the small change in mass value. The uncertainty band also shows that under some combinations of $^{84}$Ga and $^{85}$Ga masses within their corresponding error bars the peak at $A=84$ is severely under–produced.

![Figure 1](image)

*Figure 1. (Upper plot) Final abundances averaged over calculations with $Y_e = 0.35–0.38$ compared to the solar r–process abundance, with uncertainty shown as a gray band. The pink band shows the 3σ change in calculated production, a result of the variation of the masses of $^{84,85}$Ga. The red line shows the resulted abundances using the central experimental value whereas the blue line the abundances when the AME$_{16}$ extrapolated mass values were used. (Lower plot) Change, in percentage, of the abundance pattern as a result of using the mass values from [1] compared to the extrapolations given in the AME$_{16}$. Figure adapted from [1]*

Similarly, in Fig. 2 we compare our results with the solar r–process for two separate cases; in the first case, we use data for β–delayed neutron emission branches from NUBASE16. In the second case, we use the updated data for $^{82,83,84,85}$Ga as presented in Table 2 from [33,34]. The differences are more pronounced at $A=85$ where the 2 instances of our calculations differ by $\sim 18\%$. This difference can be traced back to the changed by $\sim 35 \%$ value of $P_{1n}$, where NUBASE16 value was given as $\sim 35\%$, and [34] measurement at $\sim 70 (5) \%$.

**CONCLUSIONS**

We explored the impact of a series of newly measured Ga isotopes. We demonstrate that at moderate neutron–rich conditions, realized in BNSs, r–process calculations can produce the local peaks at $A = 80$ and $A = 84$ of the solar system r–process residual. We show that changes of only a few keV in the mass of a single nucleus can lead to differences in abundances of more than 10%. In addition, we demonstrate the impact of β–delayed neutron emissions, finding changes of $\sim 18\%$ in
abundances of A=85 calculations using the NUBASE16 data, compared to NUBASE16 updated with the $^{82,83,84,85}$Ga P$_{1n}$ values according to [33,34].

![Graph showing abundances and changes](image_url)

**Figure 2.** (Upper plot) Final abundances averaged over calculations with Ye = 0.35–0.38 compared to the solar r-process abundance, with uncertainty shown as gray band. The red line corresponds to abundances when NUBASE16 was used for β–delayed neutron emission probabilities, while the blue line corresponds to abundances when the measured [33,34] β–delayed neutron emission probabilities of Ga 82–85 were used to update NUBASE16 values. (Lower plot) Change, in percentage, of the abundance pattern as a result of using the β–decay ratios and β–delayed neutron emission probabilities values from [33,34] compared to the values given in NUBASE16.

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