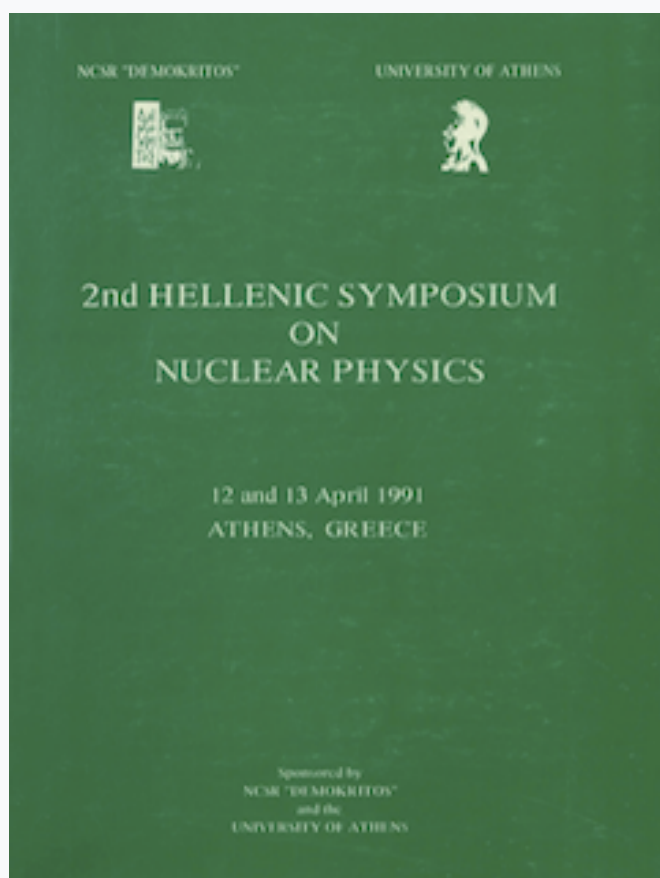


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# COMPARISON BETWEEN THE RELATIVISTIC AND NON-RELATIVISTIC TREATMENT OF THE $\Lambda$ -HYPERNUCLEI

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**Abstract:** Using as an example potentials of the form  $U_{\pm}(r) = -D_{\pm} (\cosh^2(r/R))^{-1}$ , the binding energies as well as the root mean square radii of the orbits of the  $\Lambda$  particle in hypernuclei in the ground and excited states were calculated in the relativistic and non-relativistic cases and the results are compared.

## 1. Introduction

During the works of the First Hellenic Symposium on Theoretical Nuclear Physics which took place in Thessaloniki in 1990 we have presented a study concerning  $\Lambda$ -hypernuclei in which the Dirac equation was employed. A question which remained hanging in the air at that time was the following: Do the relativistic results derived using the Dirac equation differ essentially from the non-relativistic ones in which the Schrödinger equation is used? This question we try to answer in this contribution using as an example the potential of the form

$$U(r) = -D (\cosh^2(r/R))^{-1} \quad (1)$$

The reason behind the choice of this potential is that it was used by Grypeos, Lalazissis and Massen<sup>3-5)</sup> in the non-relativistic study of the  $\Lambda$ -hypernuclei and so the comparison with the relativistic case was much easier.

## 2. Numerical results

Using the formalism outlined in ref(1) and applying a least-squares fitting procedure we have found in the relativistic case assuming that the  $\Lambda$ -nucleus potential is made up of the components

$$U_{\pm}(r) = -D_{\pm}(\cosh^2(r/R))^{-1} \quad (2)$$

that the potential parameters are

$$D_+ = 39.69 \text{ MeV}, D_- = 201.59 \text{ MeV}, r_0 = 0.984 \text{ fm.}$$

(where the extra decimals are used for the sake of comparison). We notice that the value of  $D_+$  is very close to the value of the well depth  $D$  in the non-relativistic case while the values of  $r_0$  in both cases are almost the same. Using the potential parameters given above we have calculated the binding energies of the ground state  $1s$  and of the excited states  $1p_{3/2}$  and  $1p_{1/2}$  in the relativistic case and also the binding energies of the  $1s$  and  $1p$  states in the non-relativistic case for a number of  $\Lambda$ -hypernuclei and the results obtained in both cases are given and compared in table 1.

Next using the Dirac radial wavefunctions  $G(r), F(r)$

we have calculated numerically the root mean square radii of the orbits of the  $\Lambda$  particle in the  $\Lambda$ -hypernuclei in the ground state  $1s$  and in the excited states  $1p_{3/2}$  and  $1p_{1/2}$  with the help of the formula

$$\langle r_{\Lambda}^2 \rangle^{1/2} = \left( \frac{\int_0^{\infty} r^2 (G^2(r) + F^2(r)) dr}{\int_0^{\infty} (G^2(r) + F^2(r)) dr} \right)^{1/2} \quad (3)$$

Also using the radial wave functions  $\psi(r)$  of the Schrödinger equation we have calculated numerically the root mean square radii of the orbits of the  $\Lambda$  particle in various hypernuclei in the states  $1s$ ,  $1p$  using the formula

$$\langle r_{\Lambda}^2 \rangle^{1/2} = \left( \int_0^{\infty} \psi^*(r) r^2 \psi(r) dr \right)^{1/2} \quad (4)$$

where the wavefunctions are considered normalized. The results obtained in both cases are given and compared in table 2.

### 3. Discussion

Our aim in this contribution was the comparison between the relativistic and non-relativistic results obtained in a phenomenological treatment of  $\Lambda$  hypernuclei.<sup>6)</sup> We had chosen for this comparative study the potentials (2) and (1) respectively. The quantities chosen to be compared are the binding energies of the  $\Lambda$  particle in hypernuclei as well as the root mean square radii of its orbits in them.

From tables 1 and 2 we observe that the relativistic results differ from the non-relativistic ones, as far as the binding energies are concerned, very little in the ground state namely (0.2%-0.5%) while in the excited state  $1p$  (which in the rel. case is taken as the average of the binding energies of the states  $1p_{3/2}$  and  $1p_{1/2}$ ) the difference becomes greater namely (0.7%-7%). The difference, as far as the root mean square radii are concerned is more apparent even in the ground state and is of the order of (2.1%-2.7%).

Despite the fact that the differences between the relativistic and non-relativistic treatment are not large as to make the non-relativistic calculations unreliable yet the relativistic treatment has some advantages like for instance that it incorporates the spin-orbit coupling the magnitude of which is found to be small for the  $\Lambda$ -hypernuclei an information which we cannot have with the non-relativistic treatment.

Table 1

The binding energies of various  $\Lambda$ -hypernuclei obtained relativistically and non-relativistically are given in columns II-VI and compared in VII-VIII

Relativistic				Non-Relativistic				Difference			
$U_+ = 39.7 \text{ MeV}, U_- = 201.6 \text{ MeV}$				$U = 38.93 \text{ MeV}$							
$r_0 = 0.984 \text{ fm}$				$r_0 = 0.986 \text{ fm}$							
Hyp.	$1s$ $B_\Lambda$	$1p_{3/2}$ $B_\Lambda$	$1p_{1/2}$ $B_\Lambda$	$1p$ $B_\Lambda$	$1s$ $B_\Lambda$	$1p$ $B_\Lambda$	$1p$ $B_\Lambda$	$1p$ $B_\Lambda$	$1p$ $B_\Lambda$	$1p$ $B_\Lambda$	$1p$ $B_\Lambda$
$^9\text{Be}_\Lambda$	8.66	-	-	-	8.62	-	0.37	-			
$^{13}\text{C}_\Lambda$	11.57	0.67	0.36	0.56	11.59	0.60	0.15	6.95			
$^{16}\text{O}_\Lambda$	13.12	1.96	1.57	1.83	13.15	1.92	0.27	4.68			
$^{28}\text{Si}_\Lambda$	16.91	6.00	5.58	5.86	16.95	6.04	0.25	2.98			
$^{32}\text{S}_\Lambda$	17.74	6.99	6.58	6.86	17.77	7.05	0.18	2.75			
$^{40}\text{Ca}_\Lambda$	19.07	8.64	8.25	8.51	19.09	8.72	0.13	2.39			
$^{89}\text{Y}_\Lambda$	23.27	14.26	13.95	14.15	23.23	14.36	0.16	1.44			
$^{138}\text{Ba}_\Lambda$	25.24	17.05	16.80	16.97	25.15	17.14	0.35	1.00			
$^{208}\text{Pb}_\Lambda$	26.89	19.46	19.26	19.39	26.75	19.52	0.52	0.65			

Table 2

The r.m.s. radii of the orbits of the  $\Lambda$ -particle in various hypernuclei obtained relativistically and non-relativistically are given (II-VI) and compared (VII-VIII)

Relativistic		Non-Relativistic				Difference	
Hyp	1s fm	1p <sub>3/2</sub> fm	1p <sub>1/2</sub> fm	1p fm	1s fm	1p fm	1s fm
$^9\text{Be}_\Lambda$	2.34	-	-	-	2.29	-	2.10
$^{13}\text{C}_\Lambda$	2.25	-	-	-	2.20	-	2.41
$^{16}\text{O}_\Lambda$	2.24	4.08	4.31	4.19	2.19	4.04	2.52
$^{28}\text{Si}_\Lambda$	2.27	3.50	3.56	3.53	2.21	3.43	2.62
$^{32}\text{S}_\Lambda$	2.29	3.46	3.51	3.49	2.23	3.39	2.65
$^{40}\text{Ca}_\Lambda$	2.32	3.43	3.47	3.45	2.26	3.35	2.65
$^{89}\text{Y}_\Lambda$	2.50	3.50	3.52	3.51	2.44	3.42	2.63
$^{138}\text{Ba}_\Lambda$	2.63	3.62	3.63	3.62	2.56	3.53	2.66
$^{208}\text{Pb}_\Lambda$	2.76	3.76	3.77	3.77	2.69	3.67	2.68

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